

# Cosmology of Heavy Axions: Decaying Dark Matter & Baryogenesis

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Soubhik Kumar  
UC Berkeley and LBL

w/ Josh Foster, Ben Safdi, and Yotam Soreq [2208.10504](#)  
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Loyola University, Mar 10, 2023

# Strong $CP$ Problem

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$$\mathcal{L}_{\text{SM}} \supset \frac{\alpha_3}{8\pi} \theta G\tilde{G} + Y_u \bar{Q}_L \tilde{H} u_R + Y_d \bar{Q}_L H d_R$$

- Observed  $|\bar{\theta}| = |\theta + \arg \det(Y_u Y_d)| < 10^{-10}$
- But why so small given  $\delta_{\text{CKM}} \sim \mathcal{O}(1)$ ?

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- strong  $CP$  problem**
- Solutions: ~~massless up quark~~, parity,  $CP$ , axions  
lattice results

# Axion Solution

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- Axion solution:

Peccei, Quinn '77  
Weinberg, Wilczek '78

$\bar{\theta}$  is a **dynamical** quantity

$$\mathcal{L}_{\text{SM+Axion}} \supset \frac{\alpha_3}{8\pi} \left( \bar{\theta} + \frac{a}{f_a} \right) G\tilde{G} \Rightarrow V(a) \approx -m_\pi^2 f_\pi^2 \cos \left( \frac{a}{f_a} + \bar{\theta} \right)$$

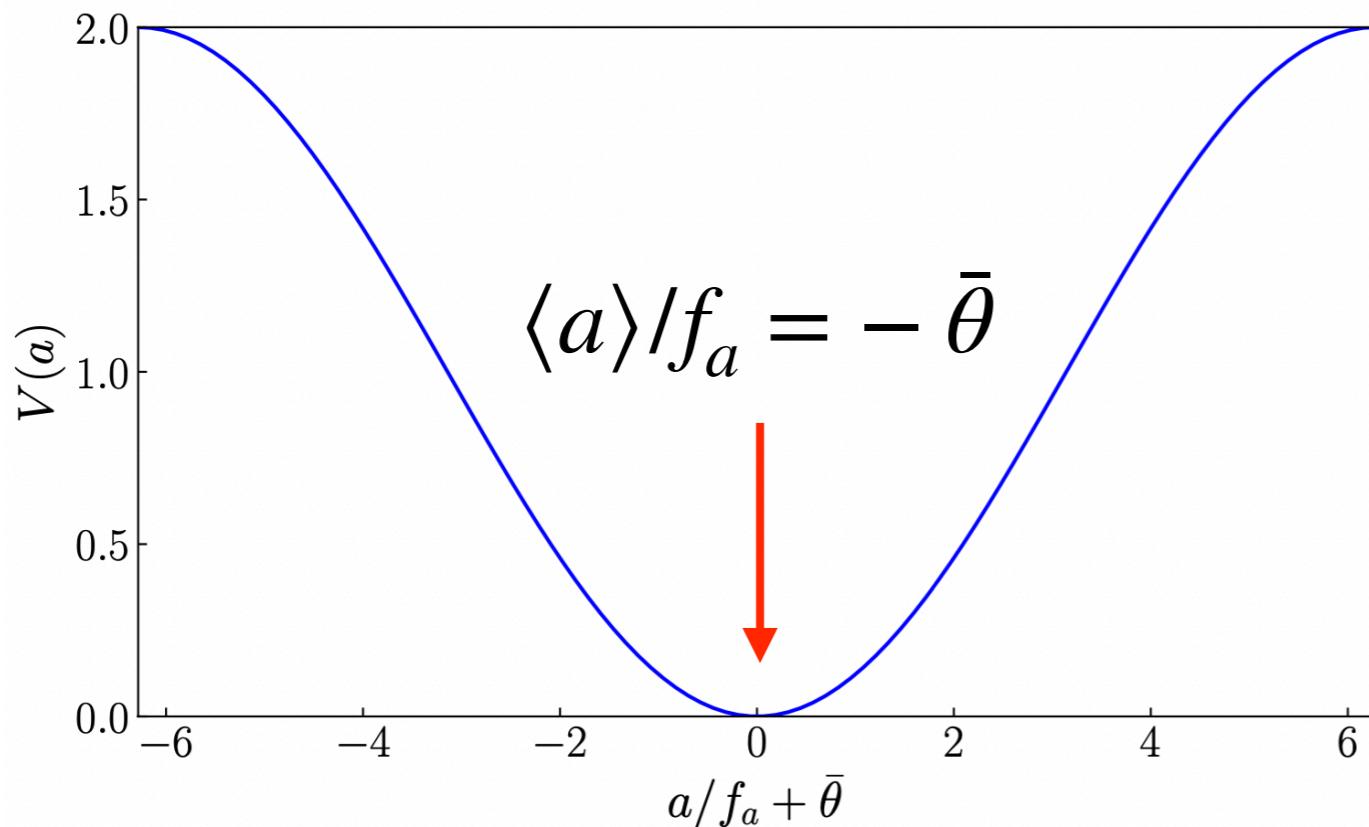
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$$\bar{\theta}_{\text{eff}} = \bar{\theta} + \langle a \rangle / f_a = 0!$$

elegant solution from  
IR EFT point of view

# Why Axions?

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- strong  $CP$  problem [QCD Axion]
- dark matter [QCD Axion or ALPs]

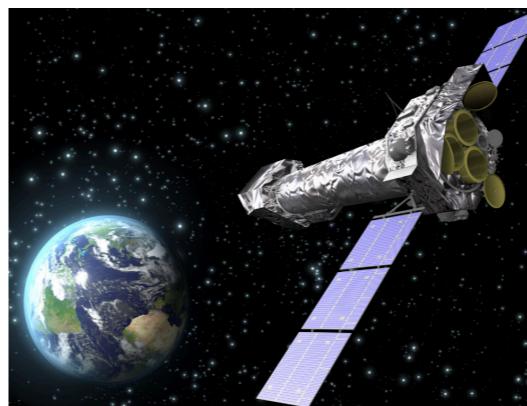
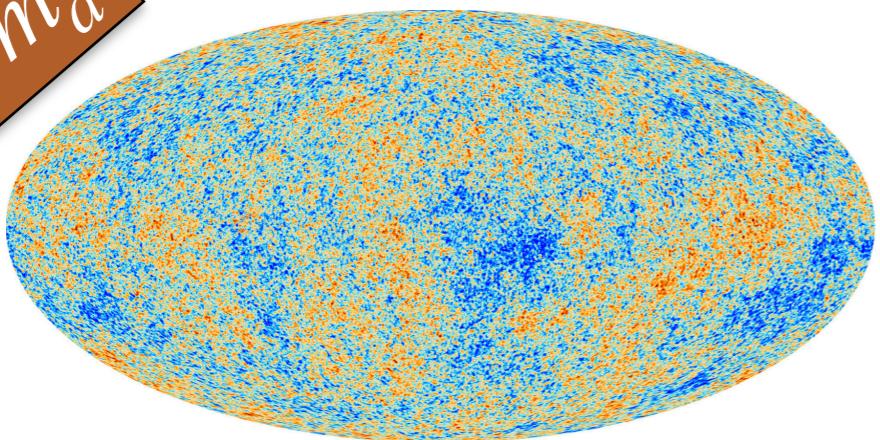
Extra-  
dim.  
Theories

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Many probes  
for small  $m_a$

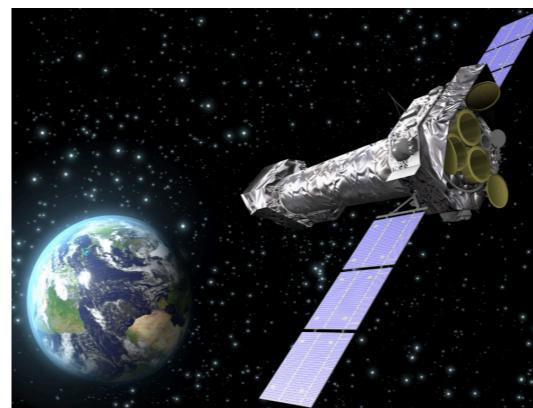
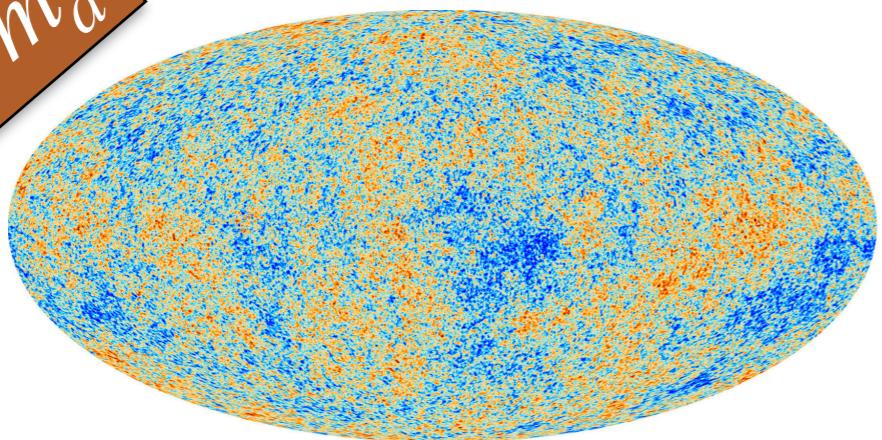


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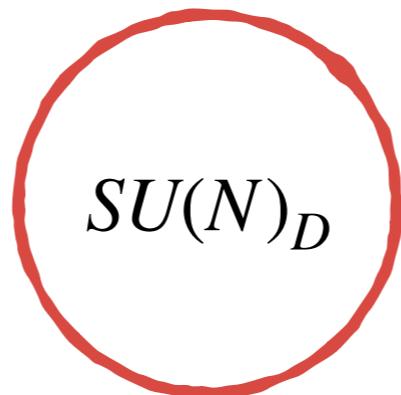
Today: **heavy axions** with  $m_a \gg \text{eV}$

- quality problem
- strong  $CP$
- dark matter

# Heavy Axions

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A toy model  
for heavy ALP<sup>†</sup>



$$\frac{\alpha_D}{8\pi f_a} a G_D \tilde{G}_D$$

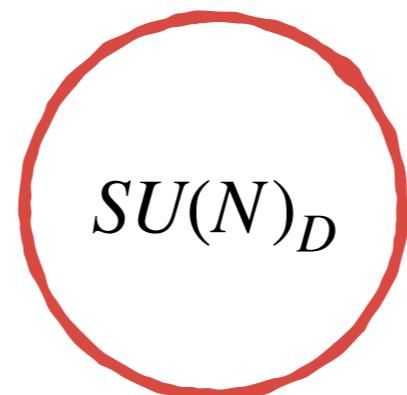
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$$m_a f_a \approx \Lambda_D^2$$

$$\Lambda_D \gg \Lambda_{\text{QCD}} \rightarrow \frac{m_a}{m_{a,\text{QCD}}} \approx \frac{\Lambda_D^2}{\Lambda_{\text{QCD}}^2} \gg 1$$

† another axion can solve  
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# How Heavy?

---

$$m_a f_a \approx \Lambda_D^2 \gg \Lambda_{\text{QCD}}^2$$

$$\Lambda_D \simeq \Lambda_{\text{UV}} \exp \left( -\frac{6\pi}{11N\alpha_{\text{UV}}} \right)$$

UV Scale                          UV Coupling

Foster, SK, Safdi, Soreq [2208.10504](#)

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$$f_a \sim \Lambda_{\text{UV}} \sim 10^{16} \text{ GeV}$$

$$\alpha_{\text{UV}} \sim 1/25$$

	$SU(2)_D$	$SU(5)_D$
$\Lambda_D$	$10^7 \text{ GeV}$	$10^{12} \text{ GeV}$
$m_a$	MeV	$10^8 \text{ GeV}$
Pheno	Dark Matter	Baryo-genesis

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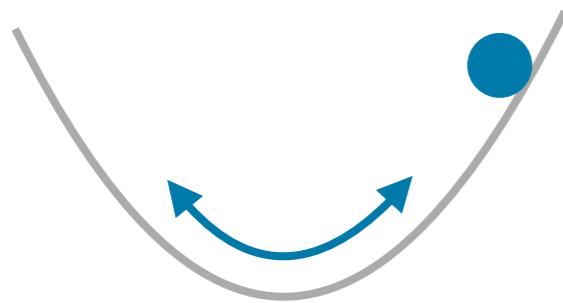


Cosmology of keV-MeV axions

Foster, SK, Safdi, Soreq [2208.10504](#)

# Heavy Axion DM

# Axion Overclosure!



Misalignment  
Production

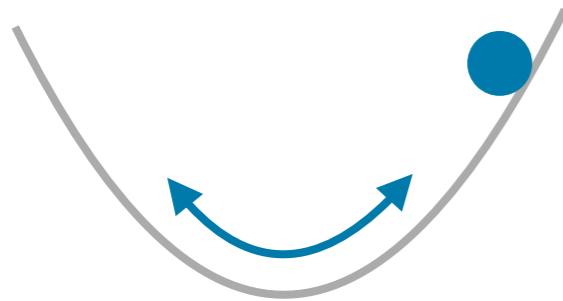
$$\boxed{m_a \sim \text{MeV}} \\ f_a \sim 10^{15} \text{ GeV}$$

$\theta_i$  = misalignment angle,  
expected  $\mathcal{O}(1)$

$$\Omega_a h^2|_{\text{RD}} \approx 0.12 \left( \frac{\theta_i f_a}{2 \times 10^{13} \text{ GeV}} \right)^2 \left( \frac{m_a}{1 \mu\text{eV}} \right)^{\frac{1}{2}} \left( \frac{90}{g_*(T_{\text{osc}})} \right)^{\frac{1}{4}}$$

$$\Omega_a \approx 10^{10} \cdot \theta_i^2 \cdot \Omega_{\text{DM}}! \:(\frown)$$

# Axion Overclosure!



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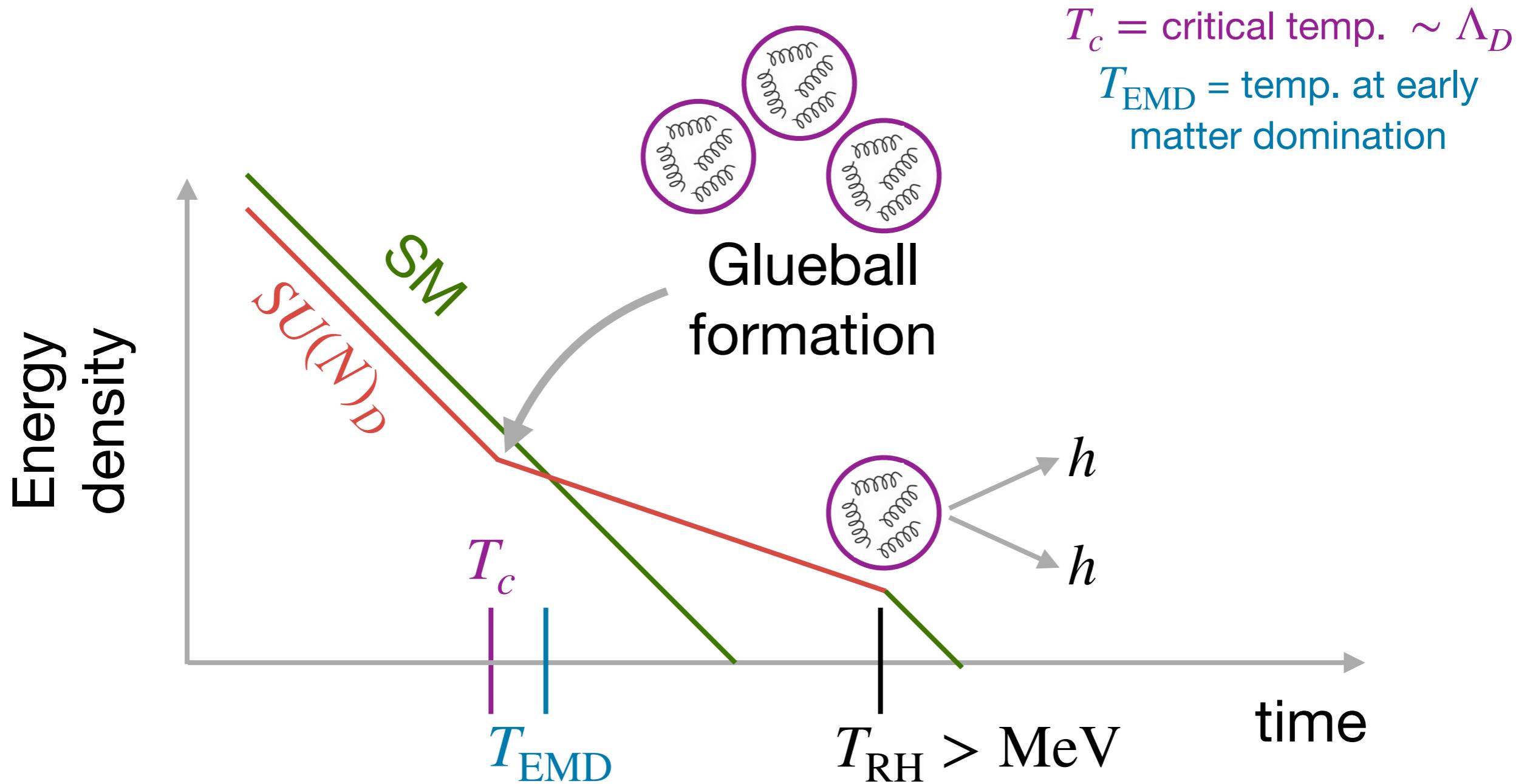
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Axion heavy due to  $SU(N)_D$

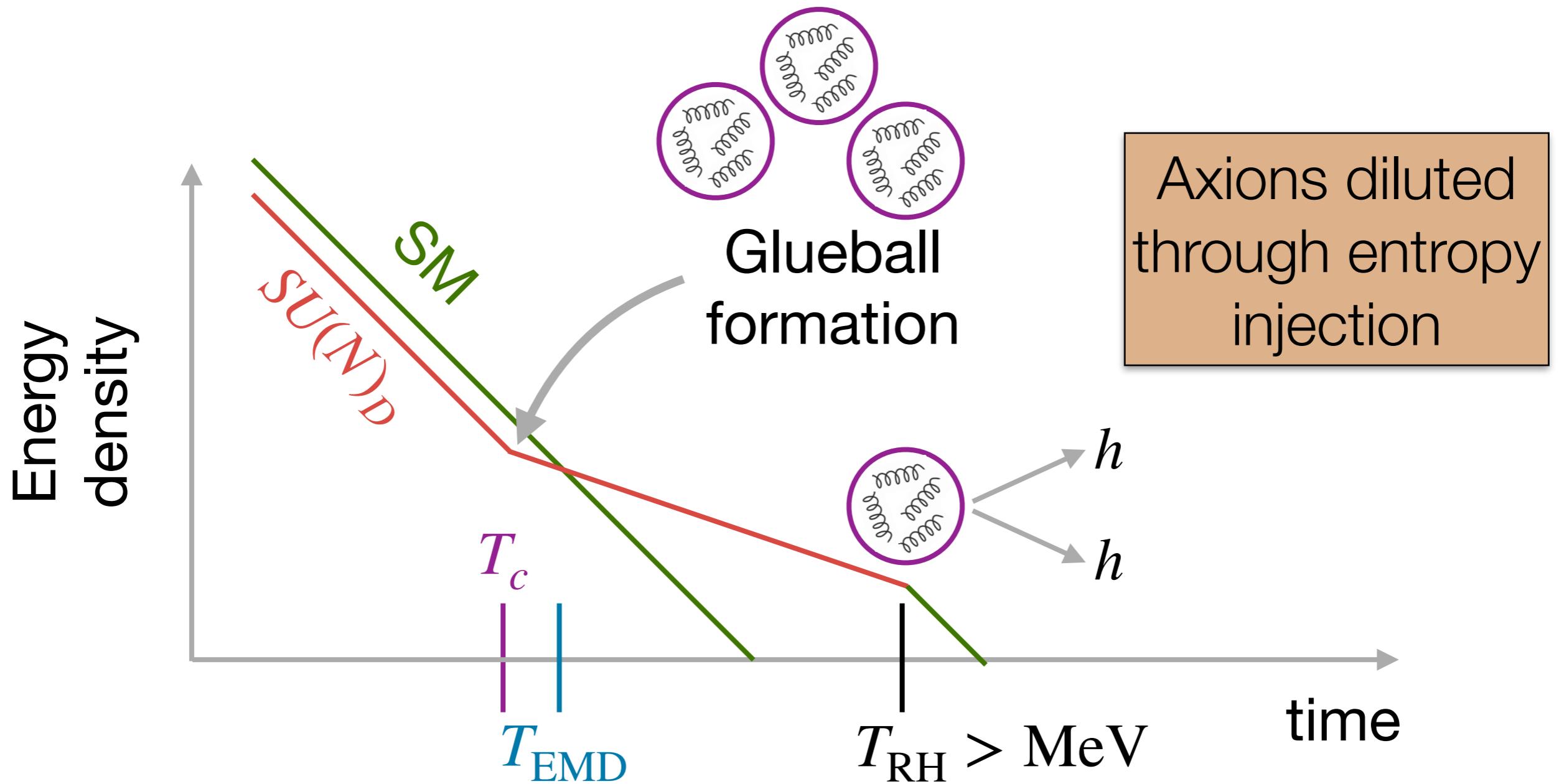


What is its cosmology?

# Cosmology of a Confining $SU(N)_D$

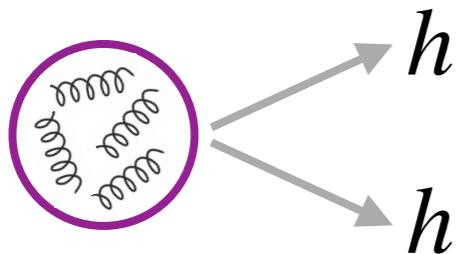


# Cosmology of a Confining $SU(N)_D$



# Dark Glueball Reheating and $\Omega_{\text{DM}}$

## Glueball Decay



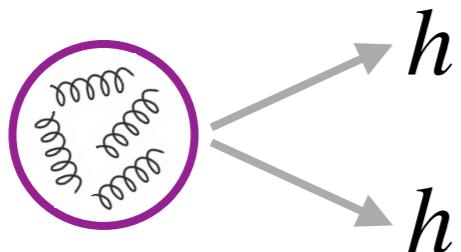
$$\frac{\alpha_D}{4\pi} G_d G_d \frac{H^\dagger H}{\Lambda^2}$$

$$\boxed{m_a \sim \text{MeV}} \\ f_a \sim 10^{16} \text{ GeV} \\ \Lambda \sim 10^{14} \text{ GeV}}$$

$$\rightarrow T_{\text{RH}} \sim 5 \text{ MeV}$$

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$$\Omega_a \approx \Omega_{\text{DM}} \left( \frac{\theta_i f_a}{10^{15} \text{ GeV}} \right)^2 \left( \frac{T_{\text{RH}}}{10 \text{ MeV}} \right)$$

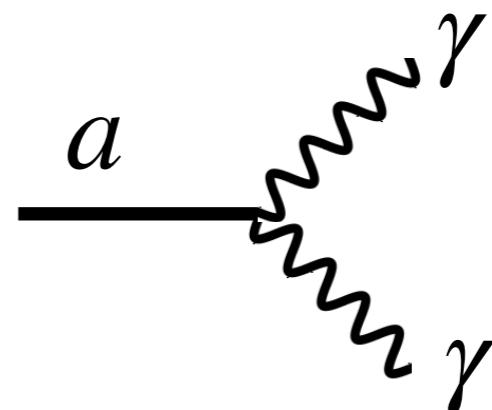


correct  $\Omega_{\text{DM}}$  w/o  
fine-tuning of  $\theta_i$ !

# Axion Lifetime

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$$\mathcal{L} \supset C_{a\gamma\gamma} \frac{\alpha_{\text{EM}}}{8\pi} \frac{a}{f_a} F^{\mu\nu} \tilde{F}_{\mu\nu}$$

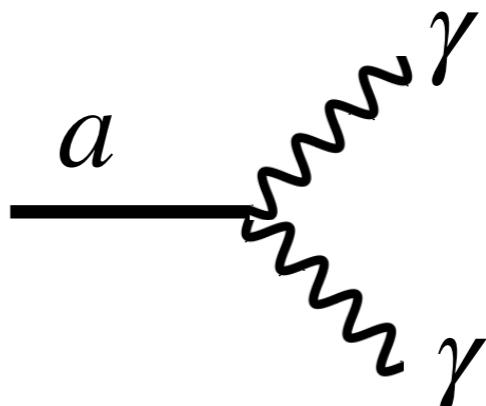


Lifetime

$$\tau_{a \rightarrow \gamma\gamma} \approx 10^{28} \text{ s} \left( \frac{0.1}{C_{a\gamma\gamma}} \right)^2 \left( \frac{0.1 \text{ MeV}}{m_a} \right)^3 \left( \frac{f_a}{10^{15} \text{ GeV}} \right)^2$$

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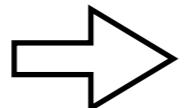


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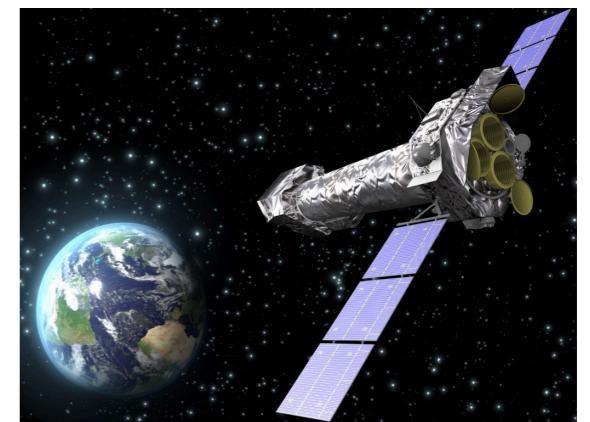
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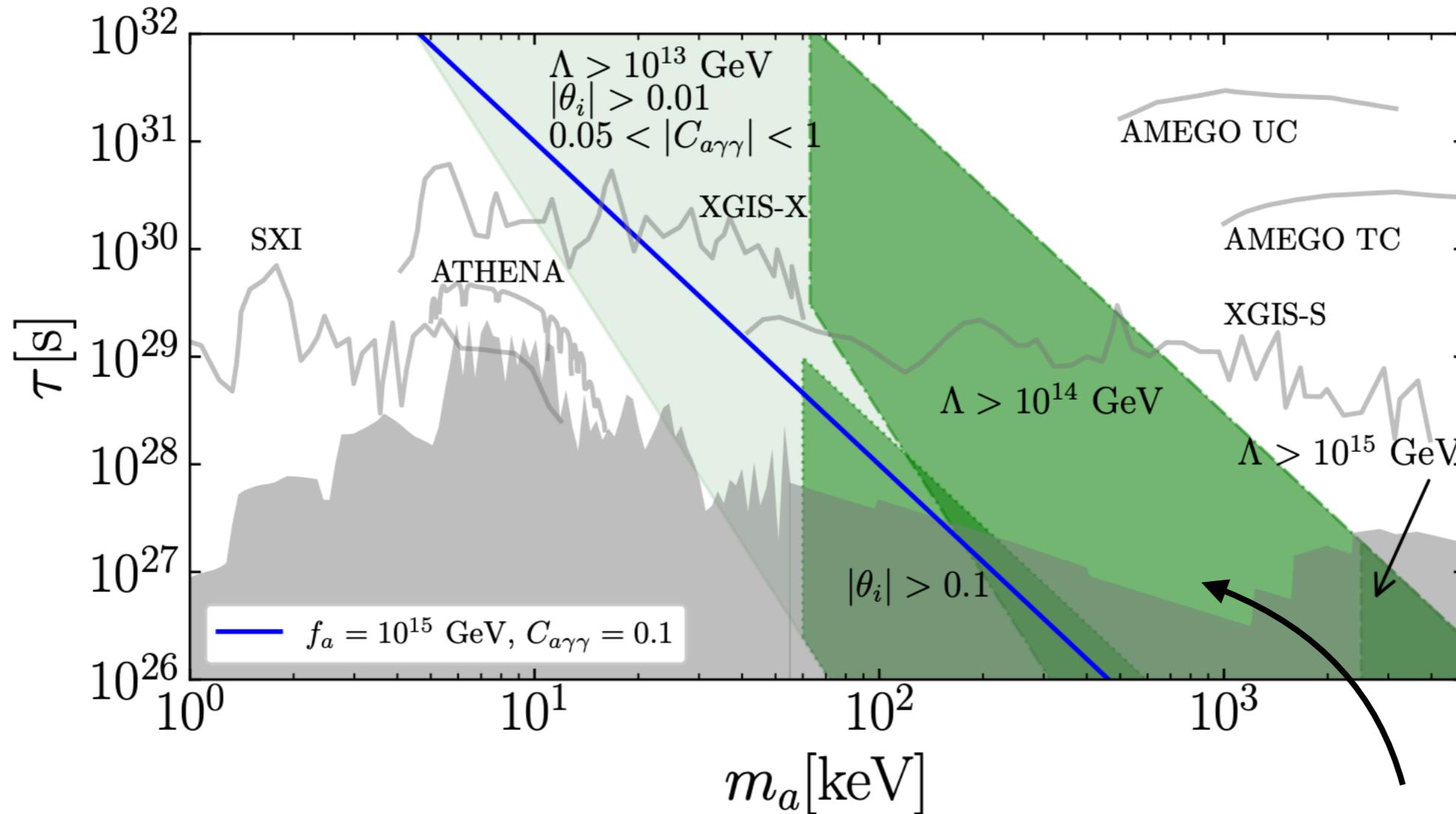
similar  $\tau$   
as decaying  
DM searches



Line signals  
at X-ray and  
 $\gamma$ -ray telescopes



# Decaying Heavy Axion DM



$SU(2)_D$  and  $T_{\text{RH}} > 5 \text{ MeV}$

Foster, SK, Safdi, Soreq [2208.10504](#)

But can we explain baryon asymmetry given low  $T_{\text{RH}}$ ?

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Cosmology of  $10^8 \text{ GeV}$  axions

Foster, SK, Safdi, Soreq [2208.10504](#)

# Spontaneous Baryogenesis

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- Fermion coupling as an example:

Cohen, Kaplan '87

$$\frac{\partial_\mu a}{f_a} \bar{\psi} \gamma^\mu \psi \supset \frac{\dot{a}}{f_a} \psi^\dagger \psi$$

spontaneous  
*CPT* breaking

- Effective “chemical potential”: fermions and anti-fermions have differing abundances

$$\Delta n_i = n_i - \bar{n}_i \approx g_i \mu_i T^2 / 6 \text{ with } \mu_i \sim \dot{a}/f_a$$

- Need *B* or *L* violation: asymmetry in equilibrium
- As *B* or *L* violation drops out of equilibrium, asymmetry also gets frozen out

# Heavy Axion Baryogenesis

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- Lepton number violation via  $(HL)^2/\Lambda_W$  with scale set by neutrino mass
- Solve Boltzmann equation for the asymmetries

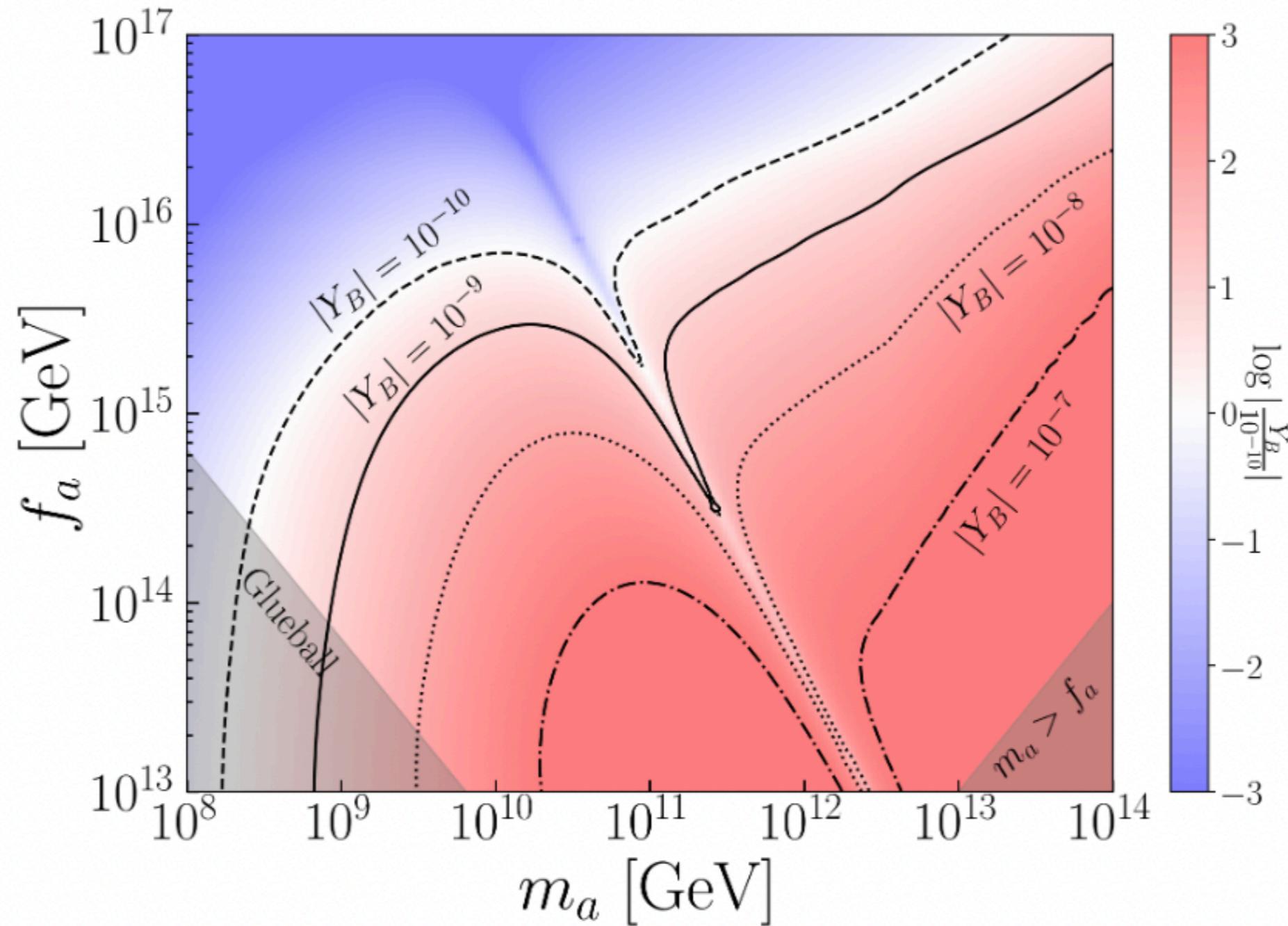
$$c_{aG} \left( \frac{\alpha_s}{8\pi f_a} aG\tilde{G} + \frac{\alpha_2}{8\pi f_a} aW\tilde{W} + \frac{\alpha_1}{8\pi f_a} aB\tilde{B} \right) + c_{af} \left( \sum_i \frac{\partial_\mu a}{f_a} J_i \right)$$

$$\frac{d}{dt} \left( \frac{\mu_i}{T} \right) = \frac{dT}{dt} \frac{1}{g_i T} \times \sum_{\alpha} \mathcal{C}_{i\alpha} \frac{\Gamma_{\alpha}}{H} \left( \sum_j \left( \frac{\mu_j}{T} \right) \mathcal{C}_{j\alpha} - n_{S\alpha} \left( \frac{\dot{a}}{f_a T} \right) \right)$$

$i = \tau, L_{12}, L_3, q_{12}, t, b, Q_{12}, Q_3, H$

$\alpha$  runs over weak sphaleron, strong sphaleron, tau Yukawa, top Yukawa, bottom Yukawa, and the Weinberg operator

# Heavy Axion Baryogenesis

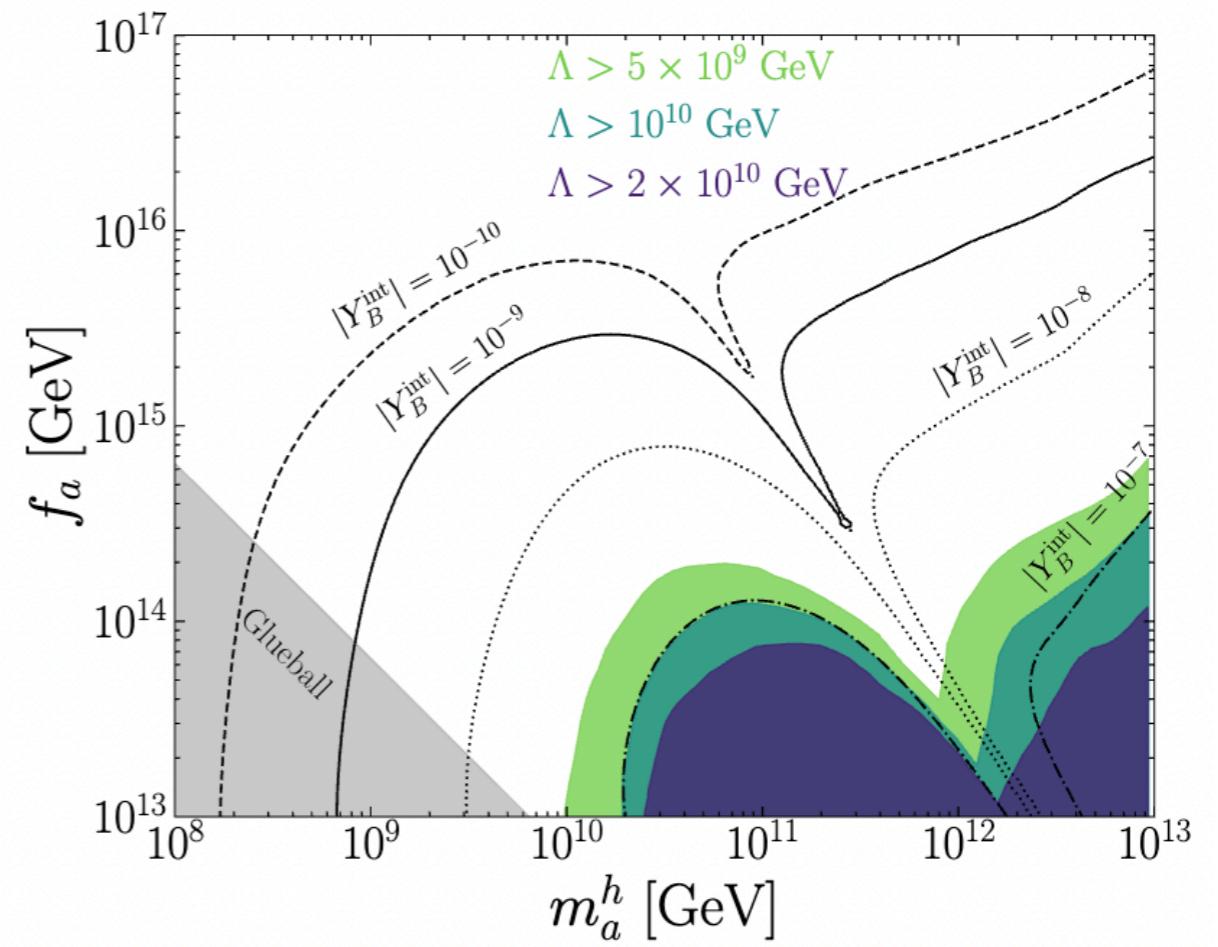
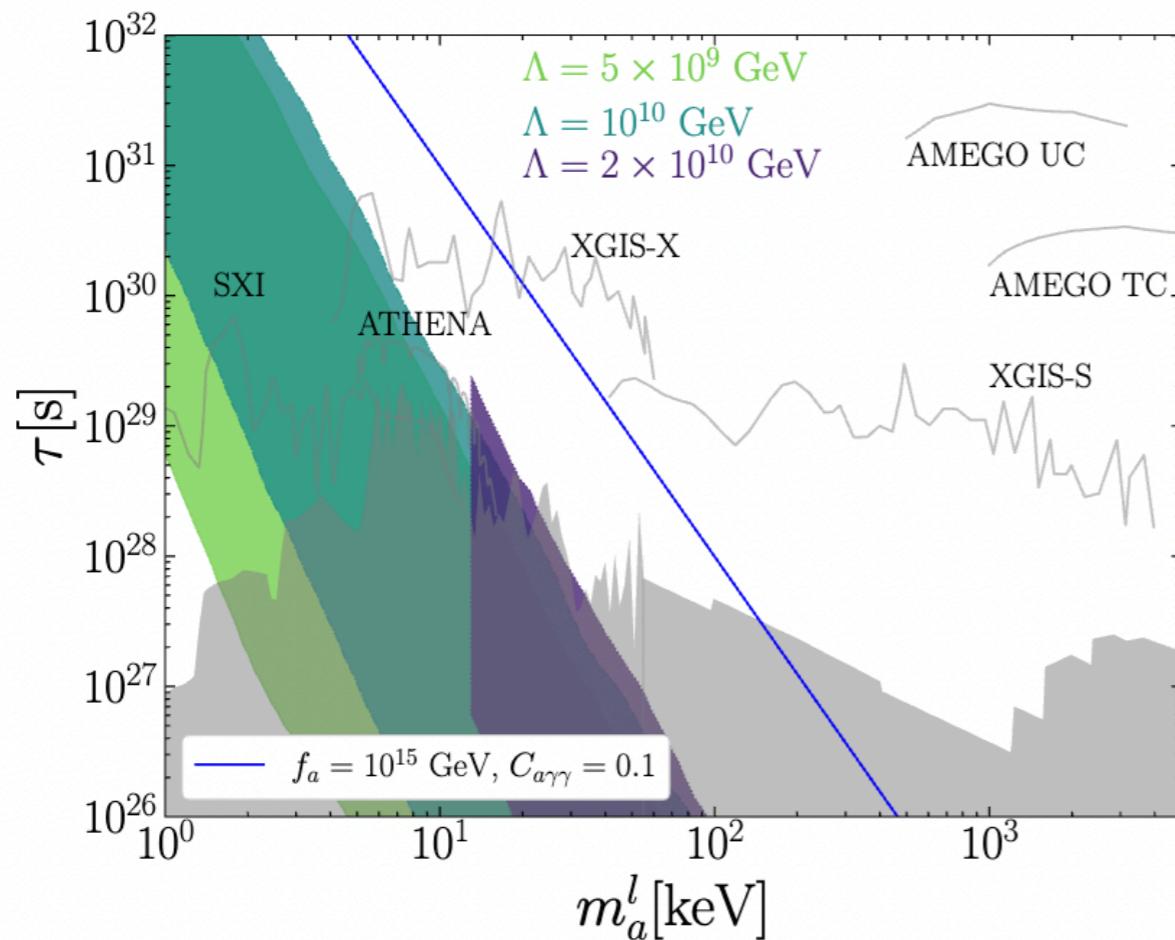


# Cosmological History

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- Two axions -  $a_h$  with  $SU(N_h)$  and  $a_l$  with  $SU(N_l)$
- $H \sim m_h$ :  $a_h$  starts to oscillate and generates initial baryon asymmetry
- Both cold and thermal  $a_h$  populations – they decay into SM
- Then  $a_l$  and glueballs of the light sector becomes important
- Eventually, entropy injections dilute both  $\Omega_{\text{DM}}$  and  $\Omega_b$

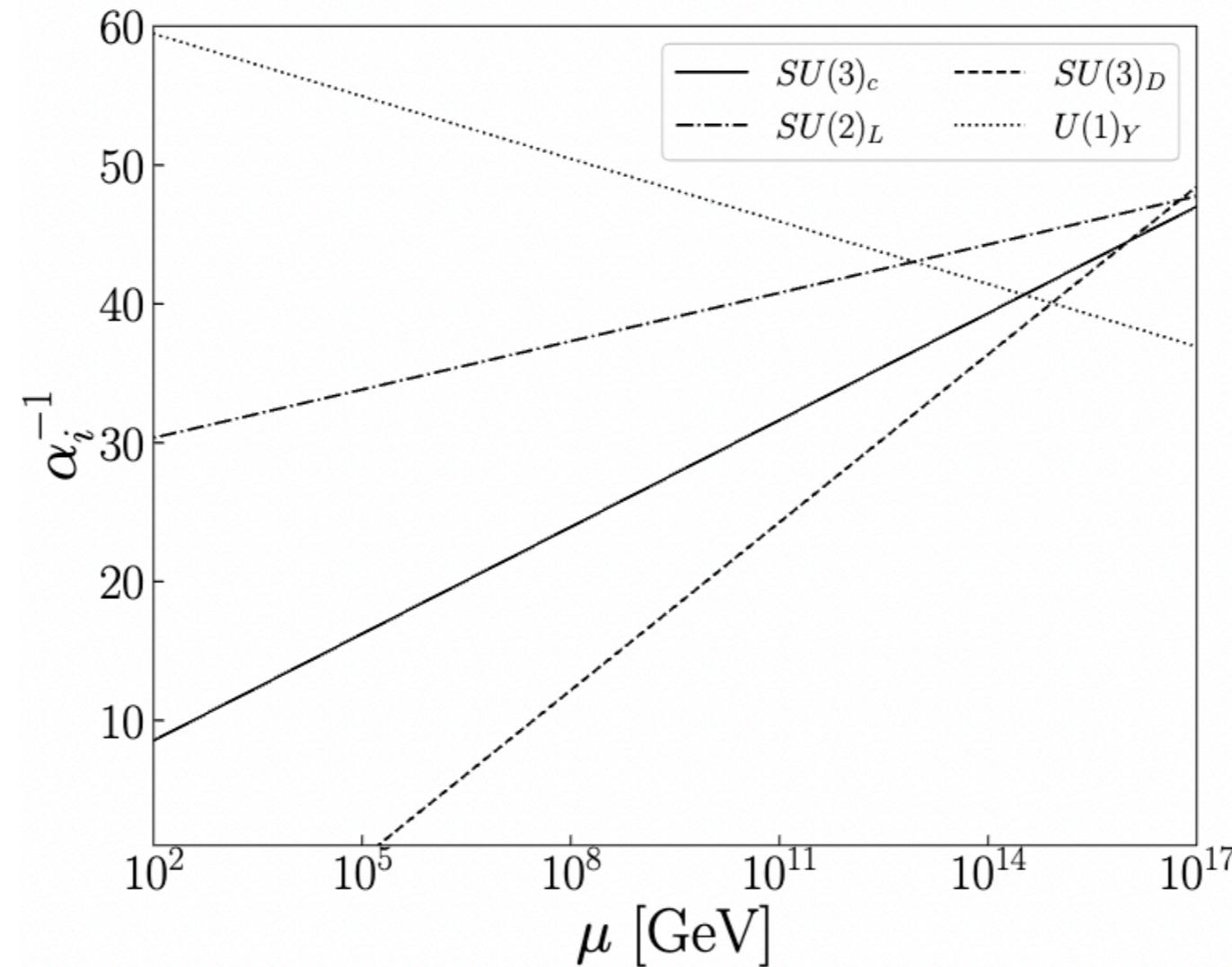
# Explaining both $\Omega_{\text{DM}}$ and $\Omega_b$



Are these  $SU(N)_D$  gauge groups  
connected to SM in a deeper way?

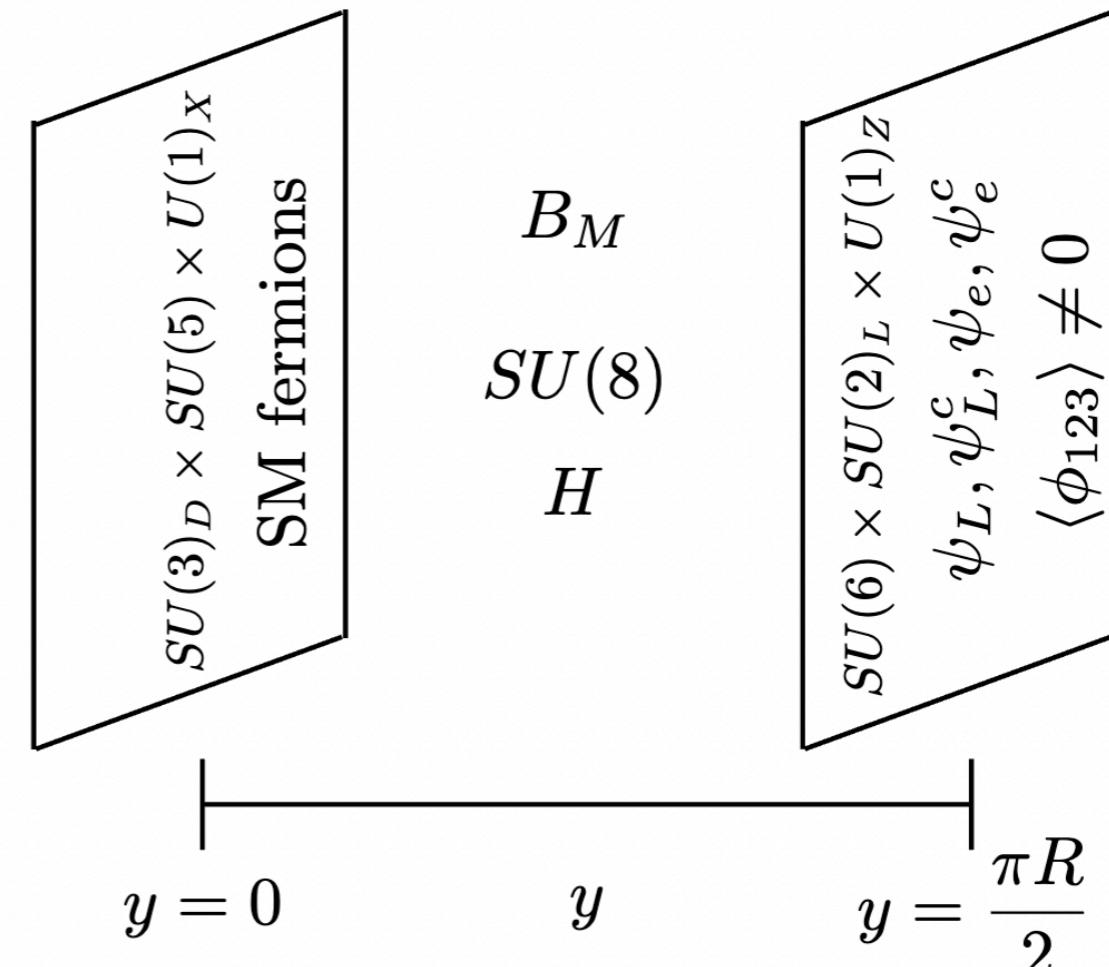
# Dark Grand Unification: Running

- Small mismatch as in non-SUSY theories, but indicates unification scale  $\sim 10^{16}$  GeV and  $\Lambda_D \sim 10^5$  GeV



# Dark Grand Unification: Construction

- Orbifold unification:  
extra dimensional  
theory with  
**boundary conditions**  
breaking gauge  
invariance
- Below the  
unification scale  
 $\sim 1/R$ :  
 $G = G_{\text{SM}} \times SU(3)_D$



Foster, SK, Safdi, Soreq [2208.10504](#)

# Axions from $U(1)$ Gauge Fields

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- Chern-Simons interaction

$$\int d^4x \int_0^{\pi R} dy \left( \frac{1}{4g_{5B}^2} B_{MN} B^{MN} + \kappa_B \epsilon^{MNPQR} B_M \text{Tr}(F_{NP} F_{QR}) \right)$$

$$\frac{\pi R}{2g_{5B}^2} (\partial_\mu B_5)^2 + 2\pi R \kappa_B B_5 G \tilde{G}$$

$$\frac{a}{32\pi^2 f_a} G \tilde{G}, \quad f_a \equiv \frac{1}{64\pi^2 \sqrt{\pi R} \kappa_B g_{5B}}$$

# Conclusion

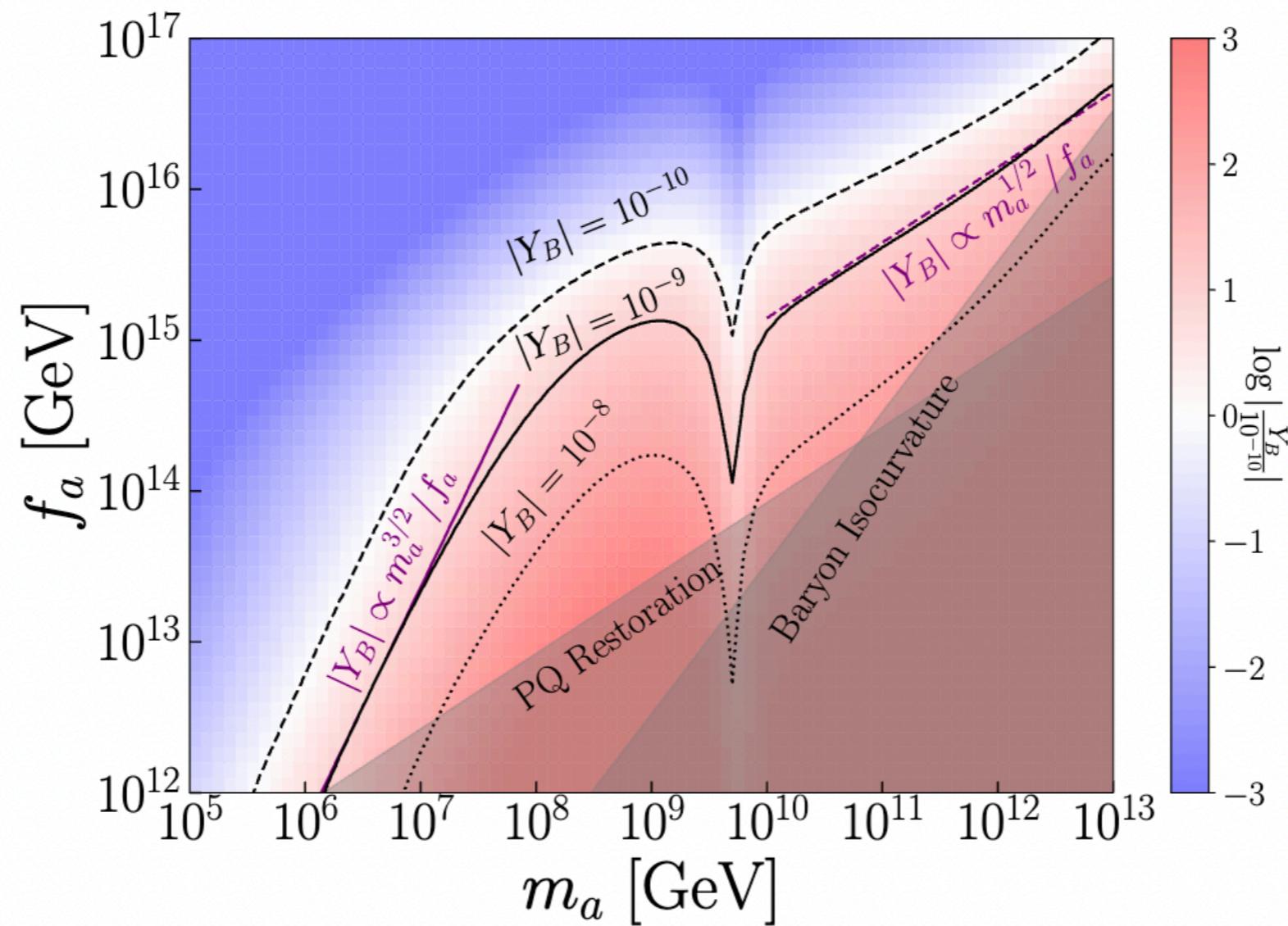
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- Heavy axions can easily arise in UV-motivated constructions
- keV-MeV scale axions can naturally explain the DM abundance
- They also predict line signals at X-ray and gamma-ray telescopes
- It is also possible to explain baryon abundance with a heavier, second axion

Thanks for your attention!

# Extra Slides

# Temperature Independent Mass



# Temperature Independent Mass

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- $\mu_{B-L} = -(\mu_\tau + 4\mu_{L_1} + 2\mu_{L_3}) + (12\mu_{q_{12}} + 3\mu_t + 3\mu_b + 12\mu_{Q_{12}} + 6\mu_{Q_3})/3$
- Coupled evolution equations

$$\dot{\rho}_a + 3H\rho_a + \frac{\rho_a}{\tau_a} = 0,$$

$$\dot{\rho}_{\text{th}} + 4H\Theta(T - m_a)\rho_{\text{th}} + 3H\Theta(m_a - T)\rho_{\text{th}} + \frac{\rho_{\text{th}}}{\tau_a} = 0,$$

$$\dot{\rho}_{\text{SM}} + 4H\rho_{\text{SM}} - \frac{\rho_a}{\tau_a} - \frac{\rho_{\text{th}}}{\tau_a} = 0,$$

$$3H^2 M_{\text{pl}}^2 = \rho_{\text{SM}} + \rho_a + \rho_{\text{th}},$$

$$\dot{\Delta n_B} + 3H\Delta n_B = 0.$$

# Axions from the UV

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- 5D Chern-Simons term:

$$\epsilon^{MNPQR} A_M \text{Tr}(F_{NP} F_{QR}) \rightarrow (a/f_a) G \tilde{G}$$

- Similar in spirit to what happens in String Theory constructions
- In fact, with multiple extra dimensions, many ways of compactifying  $\rightarrow$  many axion (-like particles) (ALPs)
- One linear combination couples to QCD  $\rightarrow$  QCD axion. Rest survive as ALPs.

# Axion Coupling

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- Linear combination coupling to the dark gauge group

$$\mathcal{L} = \sum_i \frac{\alpha_d}{8\pi} \frac{c_d^i a_i}{f_a} G_{d\mu\nu}^a \tilde{G}_{da}^{\mu\nu}$$

- Axion potential:

$$V \approx \Lambda_D^4 \left( \sum_i \frac{c_d^i a_i}{f_a} + \bar{\theta}_D \right)^2$$

- Coupling to SM gauge group:

$$\mathcal{L} = \sum_i \frac{\alpha}{8\pi} \frac{d^i a_i}{f_a} G_{\mu\nu} \tilde{G}^{\mu\nu} = \frac{\alpha}{8\pi} \frac{d_D a_D}{\tilde{f}_a} G_{\mu\nu} \tilde{G}^{\mu\nu} + \dots$$

$$d_D = \frac{\sum_i c_d^i d^i}{\sum_i (c_d^i)^2}$$

# Axion Lifetime and Dynamics

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- Into diphotons

$$\tau_{a \rightarrow \gamma\gamma} = \frac{256\pi^3}{\alpha^2 C_{a\gamma\gamma}^2} \frac{f_a^2}{m_a^3} \approx 9.6 \times 10^{27} \text{ s} \left( \frac{0.1}{C_{a\gamma\gamma}} \right)^2 \left( \frac{0.1 \text{ MeV}}{m_a} \right)^3 \left( \frac{f_a}{10^{15} \text{ GeV}} \right)^2$$

- Axion abundance

$$\Omega_a = \frac{1}{2} \frac{m_a^2 f_a^2 \theta_i^2}{\rho_c} \left( \frac{R_{\text{osc}}}{R_{\text{RH}}} \right)^3 \left( \frac{T_0}{T_{\text{RH}}} \right)^3 \frac{g_{*S}(T_0)}{g_{*S}(T_{\text{RH}})}$$

- Critical temperature

$$\frac{T_c}{\sqrt{m_a f_a}} \approx 1.6 - \frac{0.8}{N_c^2}$$

# Temperature Dependence

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- T dependent mass:  $m_a(T) \approx \begin{cases} m_a \left( \frac{T_c}{T} \right)^b, & T > T_c \\ m_a, & T \leq T_c, \end{cases}$

$$\frac{T_{\text{osc}}}{T_c} \approx \left( \frac{3\sqrt{10}}{q_T \pi \sqrt{g_*} \left( 1.6 - \frac{0.8}{N_c^2} \right)^2} \frac{M_{\text{pl}}}{f_a} \right)^{\frac{1}{2+b}} \approx \begin{cases} 5.5 \left( \frac{10^{15} \text{ GeV}}{f_a} \right)^{0.27}, & N_c = 2 \\ 2.6 \left( \frac{10^{15} \text{ GeV}}{f_a} \right)^{0.18}, & N_c = 3 \end{cases}$$

$$\Omega_a h^2 \approx 0.12 \theta_i^2 \begin{cases} \left( \frac{f_a}{10^{13} \text{ GeV}} \right)^{1.27} \left( \frac{T_{\text{RH}}}{10 \text{ MeV}} \right), & N_c = 2 \\ \left( \frac{f_a}{4.3 \cdot 10^{12} \text{ GeV}} \right)^{1.18} \left( \frac{T_{\text{RH}}}{10 \text{ MeV}} \right), & N_c = 3. \end{cases}$$